VORTEX RINGS IN HEAVY-ION COLLISIONS

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²The Ohio State University, Columbus, Ohio, USA

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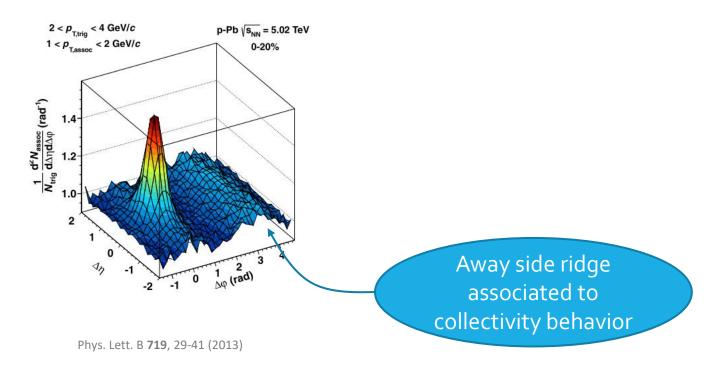
 ${}^4Wayne\ State\ University,\ Detroit,\ Michigan,\ USA\ |\ {}^5Brook haven\ National\ Laboratory,\ Upton,\ USA\ |\ {}^5Brook haven\ National\ N$

Small systems: pA collisions

- It was not expected hydro-like signals at smaller system
 - Shorter lived:
 - It was not expected it would have time for thermalization
 - It was not expected it would have time to develop collectivity behavior
 - They also have less degrees of freedom, leading many to think a thermodynamics would have limited applicability

pA collisions: The smallest fluid droplet?

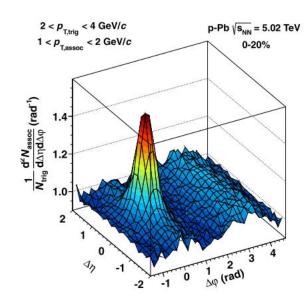
 Presence of away side in 2-particle correlation at LHC

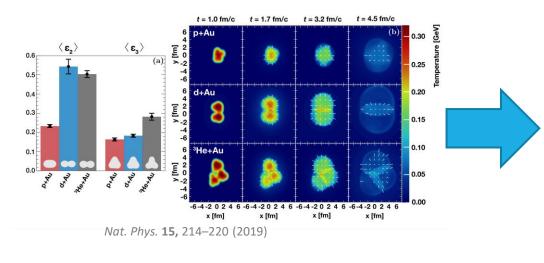


pA collisions: The smallest fluid droplet?

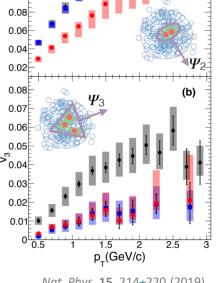
• Presence of away side in 2-particle correlation at LHC

 Elliptic flow also observed in p+Au, d+Au and 3He+Au





Phys. Lett. B 719, 29-41 (2013)



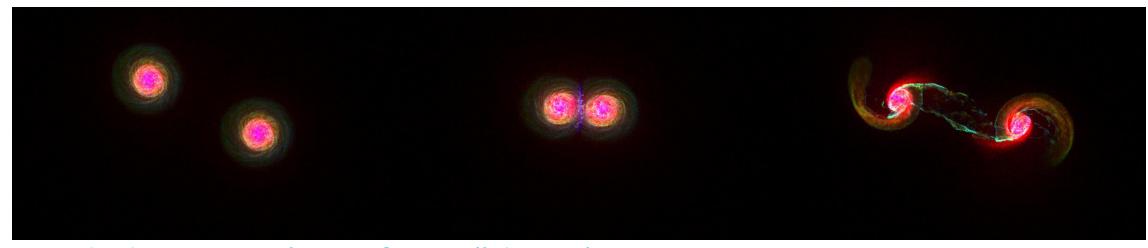
\s_{NN} = 200 GeV 0-5%

Nat. Phys. 15, 214220 (2019)

ARETHERE OTHER PHENOMENA HYDRODYNAMICS PREDICTS?

And does it happen in the QGP?

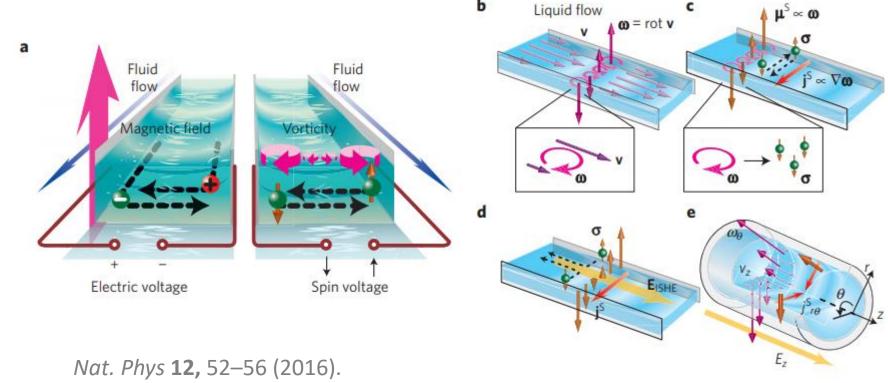
From galaxies...

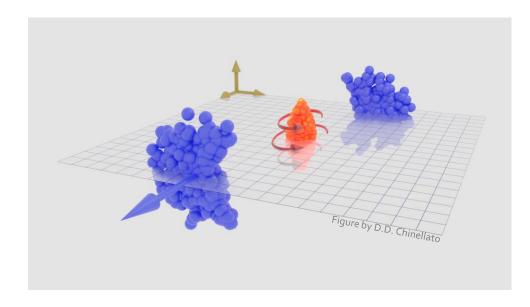


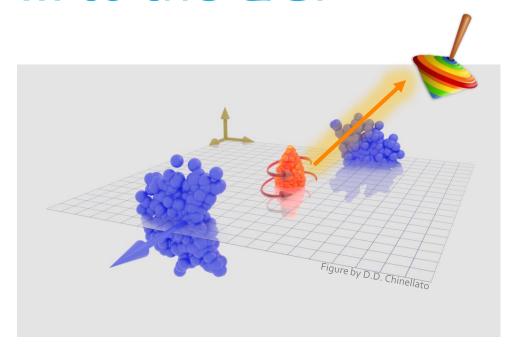
• Hydrodynamic simulation of two colliding galaxies

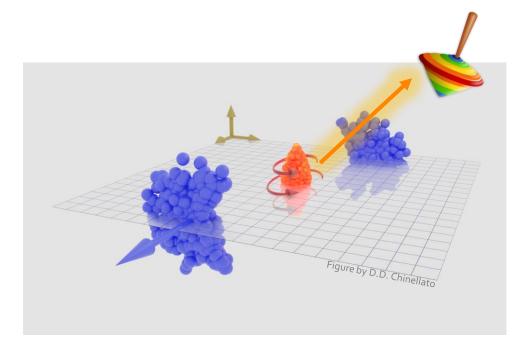
https://www.cct.lsu.edu/~werner/SphGal/

From galaxies... to the lab...

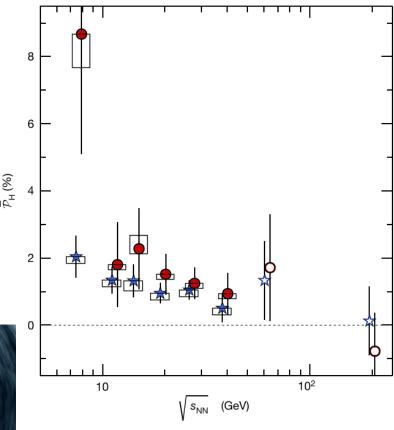






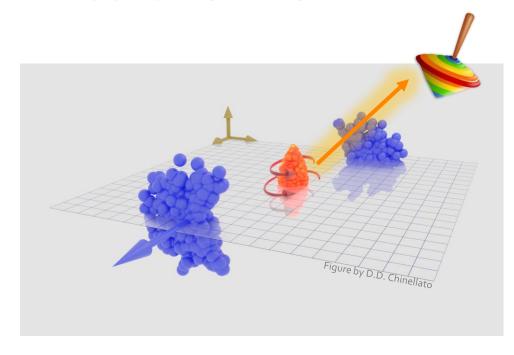


$$S^{\mu} = -\frac{1}{8m} \varepsilon^{\mu\rho\sigma\tau} p_{\tau} \frac{\int d\Sigma_{\lambda} p^{\lambda} (1 - n_{F}) \omega_{\rho\sigma}}{\int d\Sigma_{\lambda} p^{\lambda} n_{F}}$$
$$n_{f} = \frac{1}{e^{\beta^{\mu}(x)p_{\mu} - \xi(x)}}$$



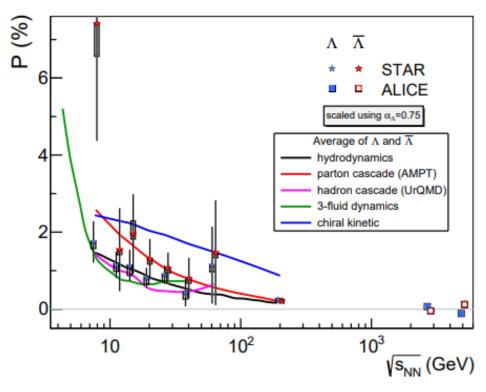
Nature **548**, 62–65 (2017)

SUBATOMIC SWIRLS



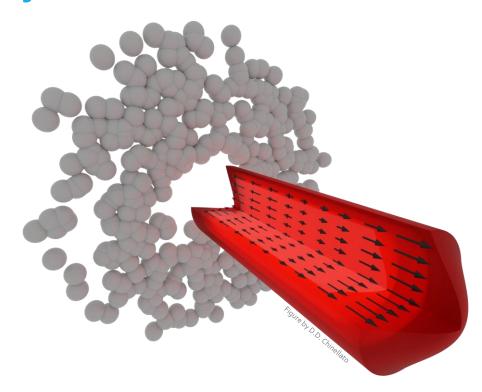
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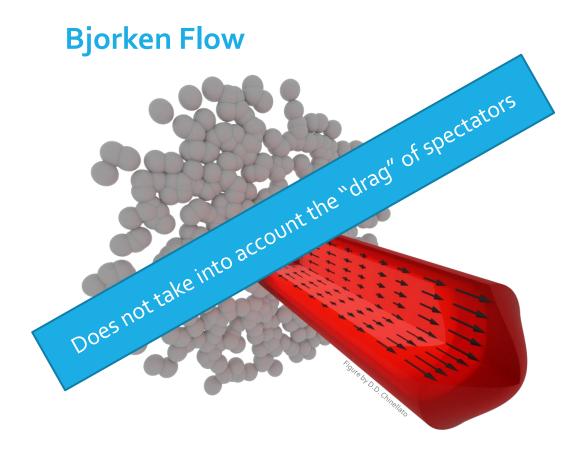
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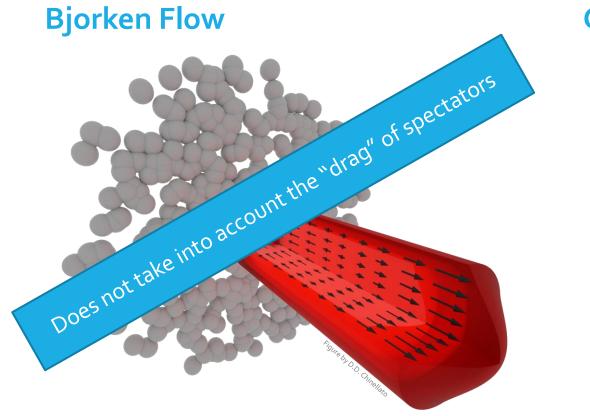


Ann. Rev. Nucl. Part. Sci. 70, 395-423 (2020)

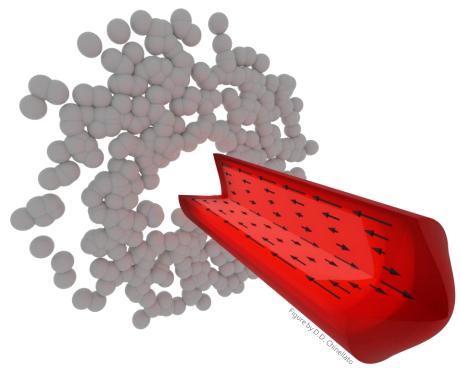
Bjorken Flow

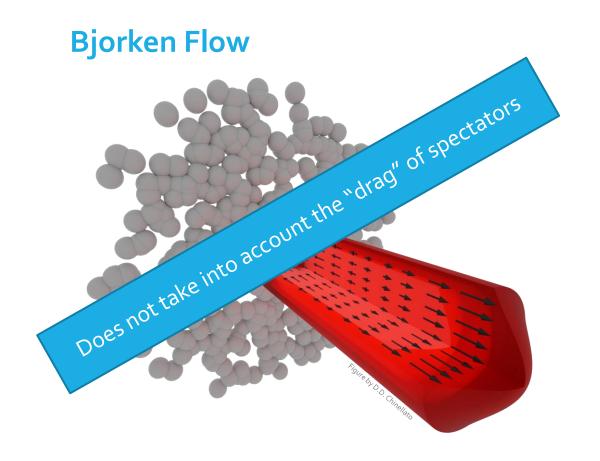




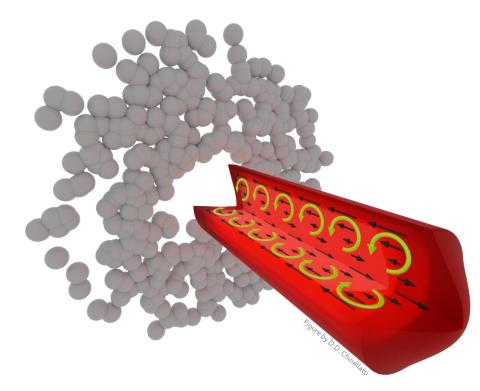


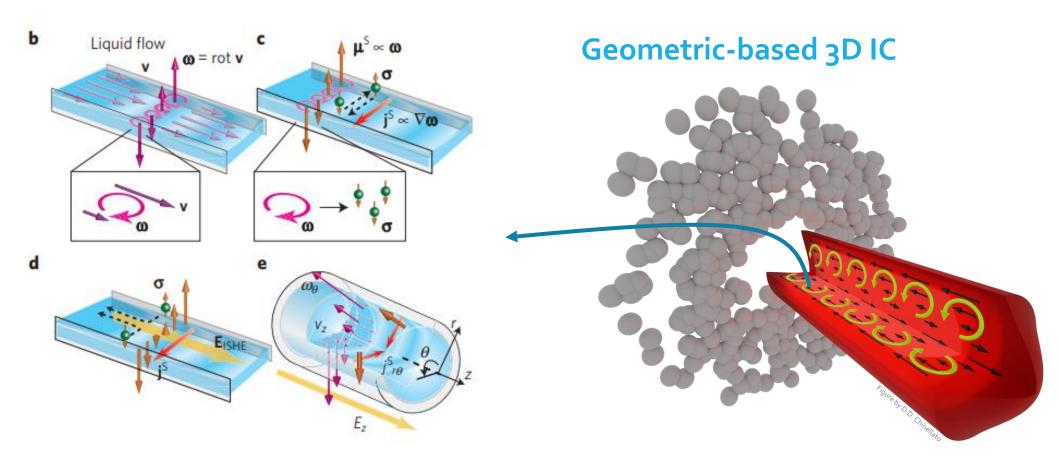






Geometric-based 3D IC





Geometric-based 3D IC

arXiv:2203.15718

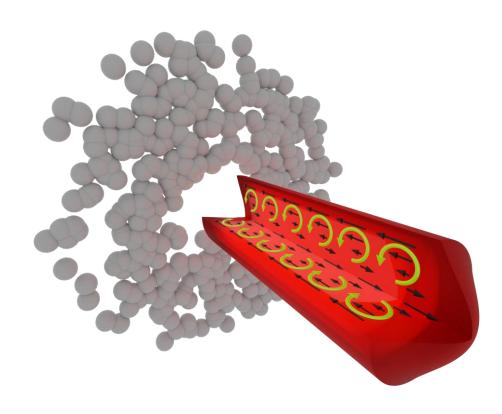
•
$$T^{\tau\tau} = e(\vec{x}_{\perp}, \eta_s) \cosh y_L(\vec{x}_{\perp})$$

•
$$T^{\tau\eta} = \frac{1}{\tau_0} e(\vec{x}_\perp, \eta_S) \sinh y_L(\vec{x}_\perp)$$

•
$$y_L(\vec{x}_\perp) = f y_{CM}(\vec{x}_\perp)$$

•
$$y_{CM}(\vec{x}_{\perp}) = \operatorname{arctanh}\left[\frac{T_A(\vec{x}_{\perp}) - T_B(\vec{x}_{\perp})}{T_A(\vec{x}_{\perp}) + T_B(\vec{x}_{\perp})} \operatorname{tanh}(y_{\text{beam}})\right]$$

•
$$e(\vec{x}_{\perp}, \eta_S; y_{\text{CM}} - y_L) = N_e(x, y) \exp \left[-\frac{(|\eta_S - (y_{CM} - y_L)| - \eta_0)^2}{2\sigma_{\eta}^2} \theta(|\eta_S - (y_{CM} - y_L)| - \eta_0) \right]$$



Geometric-based 3D IC

arXiv:2203.15718

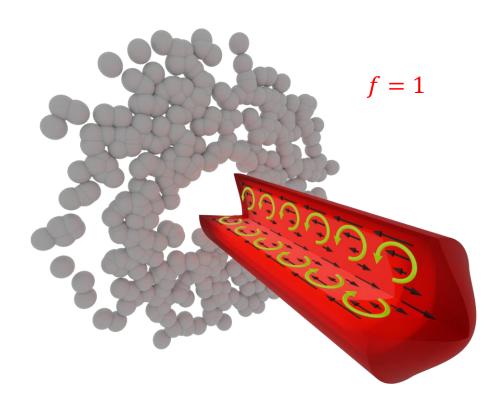
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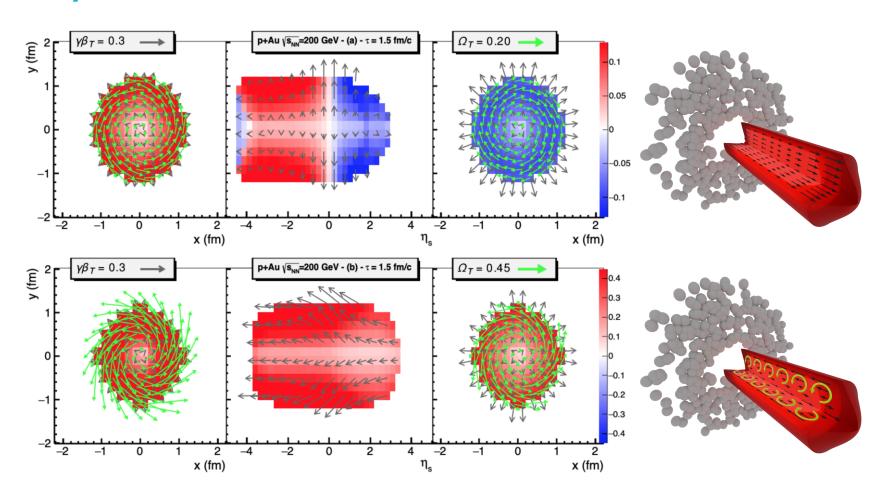
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Hydrodynamic evolution



•
$$S^{\mu} = -\frac{1}{8m} \varepsilon^{\mu\rho\sigma\tau} p_{\tau} \frac{\int d\Sigma_{\lambda} p^{\lambda} (1 - n_{F}) \omega_{\rho\sigma}}{\int d\Sigma_{\lambda} p^{\lambda} n_{F}}$$

- Three different kinds of vorticity
 - Thermal: $\omega^{\mu\nu} = \frac{1}{2} \left[\partial^{\mu} \left(\frac{u^{\nu}}{T} \right) \partial^{\nu} \left(\frac{u^{\mu}}{T} \right) \right]$
 - Kinetic: $\omega^{\mu\nu} = \frac{1}{2T} [\partial^{\mu} u^{\nu} \partial^{\nu} u^{\mu}]$
 - Temperature: $\omega^{\mu\nu} = \frac{1}{2T^2} [\partial^{\mu} (Tu^{\nu}) \partial^{\nu} (Tu^{\mu})]$

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$$S^{\mu} = -\frac{1}{8m} \varepsilon^{\mu\rho\sigma\tau} p_{\tau} \frac{\int d\Sigma_{\lambda} p^{\lambda} (1 - n_{F}) \omega_{\rho\sigma}}{\int d\Sigma_{\lambda} p^{\lambda} n_{F}}$$

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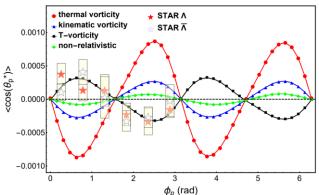
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Favored by

theory

•
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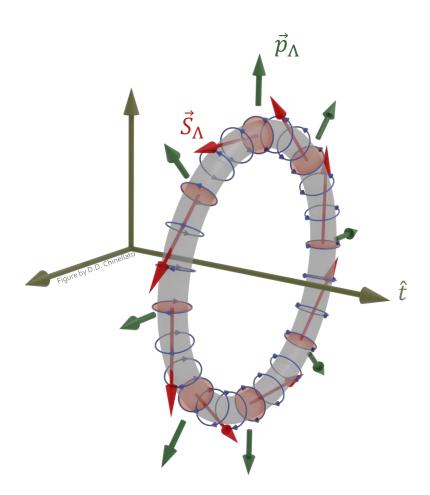


Favored by longitudinal polarization data

Phys. Rev. Research. 1, 033058 (2019). arXiv:1906.09385. I. Karpenko. arXiv:2101.04963

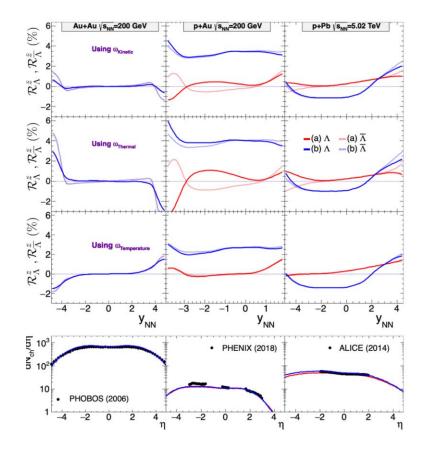
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 - Temperature: $\omega^{\mu\nu} = \frac{1}{2T^2} \left[\partial^{\mu} (Tu^{\nu}) \partial^{\nu} (Tu^{\mu}) \right]$
- Vorticity generated by smoke rings tends to cancel each other when averaged
 - Ring observable: $R_{\Lambda}^{\hat{t}} = 2 \left\langle \frac{\vec{S}_{\Lambda} \cdot (\hat{t} \times \vec{p}_{\Lambda})}{|\hat{t} \times \vec{p}_{\Lambda}|} \right\rangle_{\phi}$



Results for Smooth Initial conditions

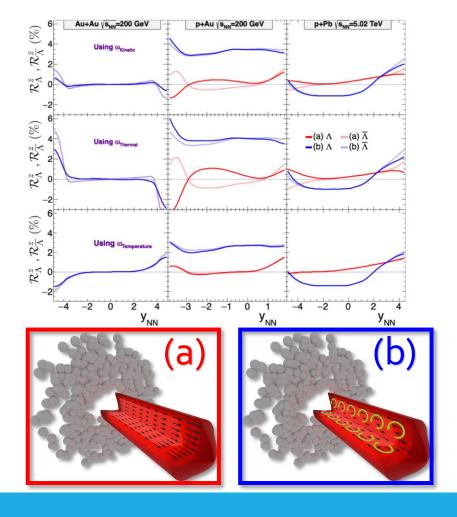
See Phys. Rev. C 104 (2021) 1, 011901



• Negligible effects on $dN/d\eta$ of charged particles

Results for Smooth Initial conditions

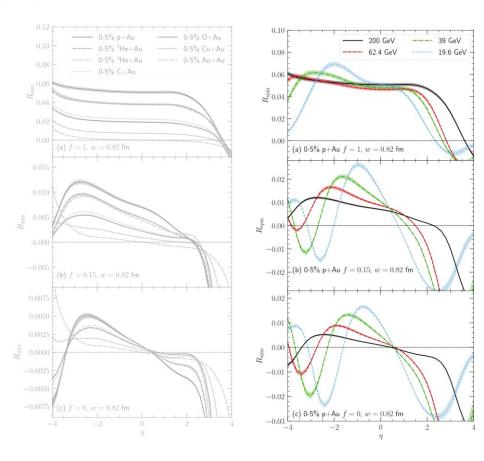
See Phys. Rev. C 104 (2021) 1, 011901



- Negligible effects on $dN/d\eta$ of charged particles
- Large ring polarization present in pA collisions at all centralities for geometric based IC
 - Effect more pronounced at RHIC than LHC
- For Bjorken-like IC:
 - Bigger signal at large rapidity

Fluctuating initial condition

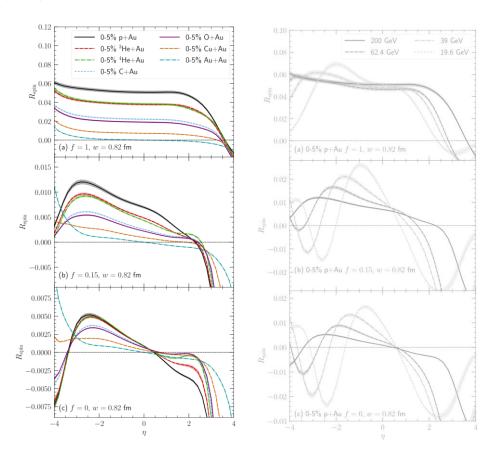
In preparation



- Collision energy (p+Au)
 - Full 3D geometric based IC
 - Signal is ~50% bigger than smooth
 - More sensitive to change of collision energy at forward rapidity
 - Bjorken based IC is sensitive to changes in energy collision at all energies

Fluctuating initial condition

In preparation



- Collision energy (p+Au)
 - Full 3D geometric based IC
 - Signal is ~50% bigger than smooth
 - More sensitive to change of collision energy at forward rapidity
 - Bjorken based IC is sensitive to changes in energy collision at all energies
- System size ($\sqrt{s_{NN}} = 200 \, \text{GeV}$)
 - The way R_{spin} changes with energy can act as a probe to the parameter f

Corrections to thermal vorticity

Symmetric shear induced polarization

•
$$S^{\mu} \rightarrow S^{\mu} + \langle A^{\mu} \rangle$$

•
$$A^{\mu} = \frac{\varepsilon^{\mu\rho\tau\sigma}}{E} p_{\tau} \xi_{\rho\lambda} \times \begin{cases} \hat{t}_{\rho} p^{\lambda} \\ \hat{u}_{\rho} p_{\perp}^{\lambda} \end{cases}$$

•
$$\xi_{\rho\lambda} = \frac{1}{2} \left[\partial_{\rho} \left(\frac{u_{\lambda}}{T} \right) + \partial_{\lambda} \left(\frac{u_{\rho}}{T} \right) \right]$$

•
$$\hat{t} = (1, 0, 0, 0)$$

•
$$p_{\perp}^{\lambda} = (\eta^{\sigma\lambda} - u^{\sigma}u^{\lambda})p_{\lambda}$$

Phys. Lett. B **820**, 136519 (2021) JHEP **07**, 188 (2021)

Spin Hall effect induced polarization

•
$$S^{\mu} \to S^{\mu} + \left\langle T \varepsilon^{\mu\nu\alpha\beta} u_{\nu} p_{\alpha} \partial_{\beta} \left(\frac{\mu_{B}}{T} \right) \right\rangle$$

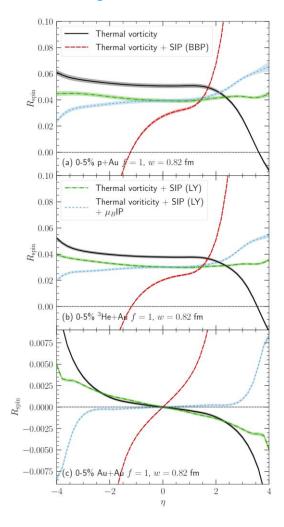
Phys. Rev. D 104, 054043 (2021)

Corrections to thermal vorticity

•
$$A^{\mu} = \frac{arepsilon^{\mu\rho\tau\sigma}}{arepsilon} p_{ au} \xi_{\rho\lambda} imes \begin{cases} \hat{t}_{
ho} p^{\lambda} & \text{BBP} \\ \hat{u}_{
ho} p_{\perp}^{\lambda} & \text{LY} \end{cases}$$

$$\bullet \left\langle T \varepsilon^{\mu\nu\alpha\beta} u_{\nu} p_{\alpha} \partial_{\beta} \left(\frac{\mu_{B}}{T} \right) \right\rangle \qquad \mu_{B} \mathsf{IP}$$

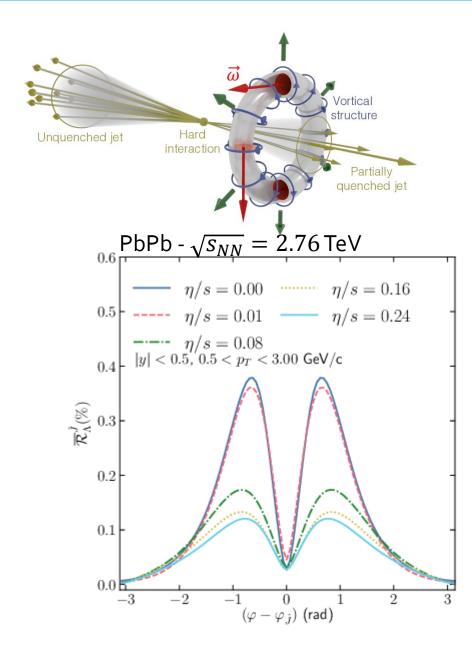
- Very sensitive to symmetric shear model
- Behavior is not intuitive
- Role of μ_B is noticeable



Applicability on jets

Phys. Lett. B 820, 136500 (2021)

- One can change the trigger direction from beam axis to jet direction
- A simple model (a hot spot with an initial velocity) has shown promising results
- A more difficult problem:
 - Predictions are very sensitive to parton-medium interaction model



Summary

- Ring observable can be used as a probe for different models in heavy-ion collisions
- Experiments should measure this now—discovery potential!
 - It is sensitive to longitudinal velocity patterns in IC
 - An energy scan can confirm a value for $f \sim 1$ if one looks at forward rapidity.
 - ullet Similarly, a system size scan can also probe values for f
 - When considering symmetric shear induced polarization, one can see large differences between the two competing models
- It may also be used to constraint parton-medium interaction models

Thanks





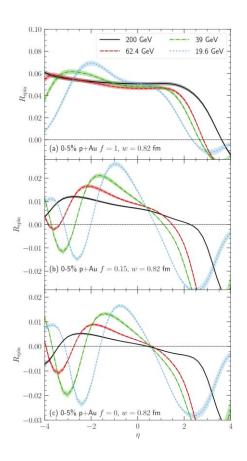


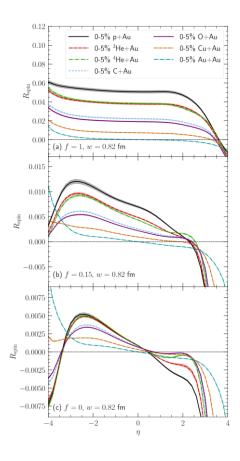
Grant # 2017/05685-2 Grant # 2018/24720-6 Grant # 2021/01670-6 Grant # 2021/01700-2

Grant # 306152/2020-7

Backup

Fluctuating initial condition





- Signal is bigger by a factor ~2
- Full 3D geometric based IC
 - Reduced sensitive to energy

Ring fluid strength

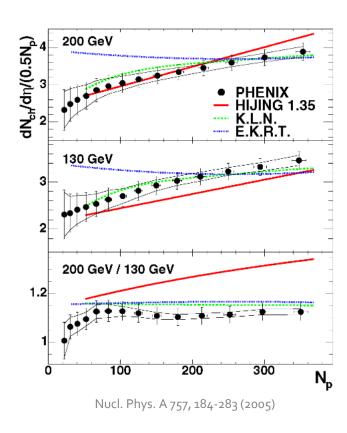
$$\bar{R}_{\text{fluid}}^{\hat{t}} = \frac{\varepsilon^{\mu\nu\rho\sigma}\Omega_{\mu}n_{\nu}\hat{t}_{\rho}u_{\sigma}}{|\varepsilon^{\mu\nu\rho\sigma}\Omega_{\mu}n_{\nu}\hat{t}_{\rho}u_{\sigma}|} \qquad \Omega_{\mu} = -\varepsilon^{\mu\nu\rho\sigma}\omega_{\rho\sigma}(u_{\nu}u^{\alpha} + C\Delta_{\nu}^{\alpha})n_{\alpha}$$

Geometric-based 3D IC

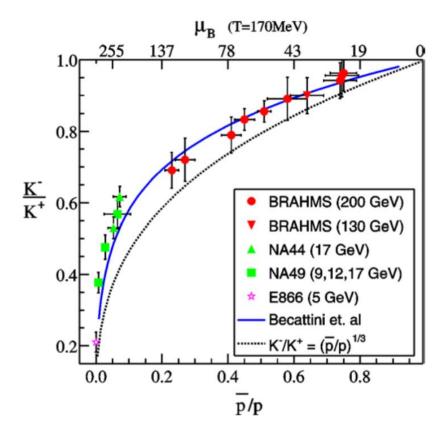
- $T^{\tau\tau} = e(\vec{x}_{\perp}, \eta_s) \cosh y_L(\vec{x}_{\perp})$
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$$e(x, y, \eta_S, y_{\text{CM}}, f) = N_e(x, y) \exp \left[-\frac{(|\eta_S - (y_{CM} - y_L)| - \eta_0)^2}{2\sigma_\eta^2} \theta(|\eta_S - (y_{CM} - y_L)| - \eta_0) \right]$$



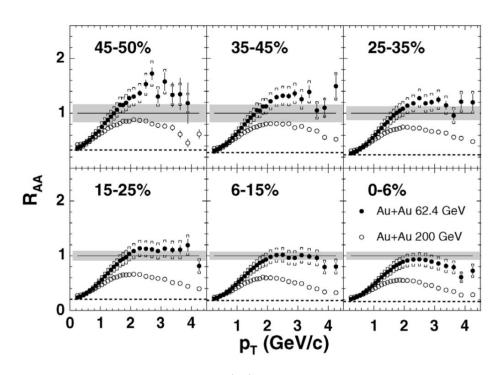


• A dense system



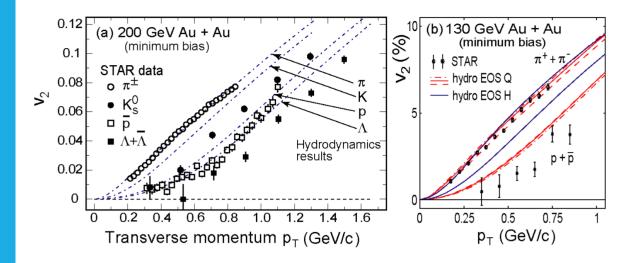
Nucl. Phys. A 757, 1-27 (2005)

- A dense system
- Hints of a thermalized system



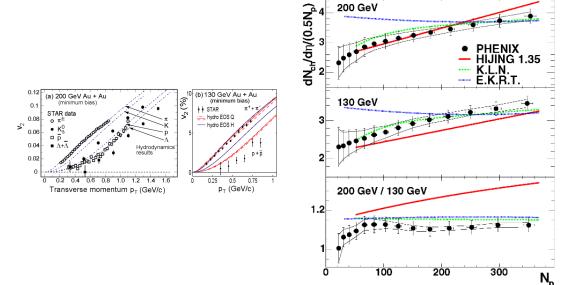
Nucl. Phys. A 757, 28-101 (2005)

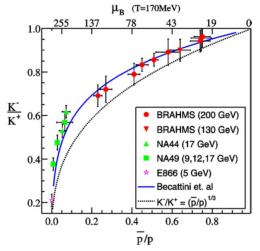
- A dense system
- Hints of a thermalized system
- A strongly interacting medium

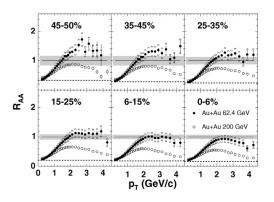


Nucl. Phys. A 757, 102-183 (2005)

- A dense system
- Hints of a thermalized system
- A strongly interacting medium
- A system which presents collectivity behavior







LOW-VISCOSITY RELATIVISTIC HYDRODYNAMICS IS ABLE TO DESCRIBE RELATIVELY WELL MOST OF THESE SIGNALS



https://www.youtube.com/watch?v=YoLKve4kofc